

Growth and magnetic characterization of Co nanoparticles obtained by femtosecond pulsed laser deposition

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We present a detailed study on the morphology and magnetic properties of Co nanostructures deposited onto oxidized Si substrates by femtosecond pulsed laser deposition. Generally, Co disks of nanometric dimensions are obtained just above the ablation threshold, with a size distribution characterized by an increasingly larger number of disks as their size diminishes, and with a maximum disk size that depends on the laser power density. In Au/Co/Au structures, in-plane magnetic anisotropy is observed in all cases, with no indication of superparamagnetism regardless of the amount of material or the laser power density. Magnetic force microscopy observations show coexistence of single-domain and vortex states for the magnetic domain structure of the disks. Superconducting quantum interference device magnetometry and x-ray magnetic circular dichroism measurements point to saturation magnetization values lower than the bulk, probably due to partial oxidation of the Co resulting from incomplete coverage by the Au capping layer.

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I. INTRODUCTION

Pulsed-laser deposition (PLD) is a versatile technique for the deposition of ultrathin films of practically any material. Recently, the use of femtosecond (fs) lasers for PLD (fs-PLD) has become an active field of research from several different aspects. The use of ultrafast lasers for ablation of materials has principally two distinctive characteristics with respect to nanosecond (ns) PLD: first, fs lasers do not interact with the ejected material since the pulse has finished well before any ablation process of the target takes place. This enables studies of the fundamental thermodynamics and hydrodynamics of particle formation. Second, short laser pulses heat up the material to much higher temperatures than nano-

second pulses with comparable fluence because the energy is delivered before any significant thermal conduction into the target's bulk can occur.

From a theoretical point of view, different approaches have tried to identify the physical process for ablation of material from a solid target upon irradiation with fs pulses. In particular, they point to different mechanisms such as liquid phase ejection and fragmentation, homogeneous nucleation and decomposition, and spinodal decomposition.¹⁻⁵ In addition, several theoretical and experimental approaches have also revealed some insight into the ultrafast laser-pulse-target interaction. This is of particular interest for understanding the incident laser's effects on the target morphol-

ogy, the resulting plasma, and the nature of the species ejected from the target.^{6–15}

In spite of this effort, however, very little has been done on the nature and properties of the material obtained from this ablation process once it is deposited on a substrate surface, either in the form of a continuous film or as nanoparticles. Deposition of continuous thin films using fs lasers has been successful in the case of oxides or nitrides,^{16–21} with ablation in the presence of a gas in the deposition chamber that reacts or interacts to some extent with the ejected material, leading to the growth of a continuous film on the substrate.

On the other hand, several groups have recently started to study the deposition and physical properties of nanoparticles obtained by fs-PLD from a variety of targets, such as Ge,²² Permalloy,²³ Si,^{24,25} Ti,²⁶ Al,²⁷ Ni,^{28,29} or Fe.³⁰ The common feature is that different materials seem to behave roughly the same when irradiated with intensities around the threshold of plasma formation, typically 10^{12} – 10^{13} W/cm², leading to nanoparticle formation, with radii of tens of nm.³¹ Most theoretical calculations^{2,4,5,10} attribute the generation of these particles to the ejection of clusters of material of different sizes down to individual atom dimensions. Experimentally, several studies have reported a maximum in the particle size distribution at some specific particle size, with a mean particle radius between 5 and 25 nm depending on the material. As we will show below a great deal of care is necessary in characterizing the size distribution at these small dimensions in order to avoid artifacts due to the resolution limit of the particular method that is used to determine the particle size. Another important issue is the actual shape of the deposited clusters of material. They are usually described as particles that one intuitively identifies with more or less spherical entities. However, if the cluster arrives at the substrate in a liquid state, it is likely to become deformed, flattening down to a disklike shape after making contact with the substrate. This geometry could also result from plastic deformation of solid particles impinging at very high ballistic velocities.

For all these reasons, fs-PLD is an interesting emergent technique for the fabrication of magnetic nanostructures with a wide size range (that could lead to a controlled monodisperse size if, for example, a quadrupolar mass filter is attached in the deposition chamber) and exhibiting interesting geometries. However, very few works related to this matter have been published so far. Reilly *et al.*²³ compared the morphology of Ni₈₀Fe₂₀ thin films using a high-average-power free-electron laser (650 fs pulse width) and an amplified Ti:sapphire system (150 fs pulse width), studying the different morphologies, crystalline nature, and magnetization curves. Ausanio *et al.*²⁸ and Amoroso *et al.*³⁰ also reported the synthesis, structure, morphology, and magnetic characterization of Ni and Fe nanoparticles or, as they call them, nanocrystal films. Very recently, Liu *et al.*²⁹ reported also on the generation of Ni nanoparticles by ultrafast PLD, but no magnetic characterization was included. In addition, the influence of surface-plasmon resonances in the magneto-optical (magnetoplasmonic) activity of noble-metal/ferromagnet/noble-metal continuous films and nanodiscs has recently been highlighted,^{32–35} and there is a growing interest in exploring similar phenomena in systems fabricated using innovative methods.

Considering the recent interest in novel experimental techniques to obtain nanostructured magnetic materials, we have carried out a complete study focusing our attention on the fabrication and characterization of Co nanoparticles obtained by fs-PLD. Here we report on a comprehensive investigation of the fabrication, size distribution, morphology, structure, and magnetic properties of Co nanoparticles deposited by fs-PLD.

II. EXPERIMENTAL APPROACH AND SAMPLE DEPOSITION

PLD was carried out in an ultrahigh-vacuum (UHV) system previously described²² using a Ti:sapphire laser (Clark MXR CPA-2001) with a pulse width of 120 fs, a wavelength of 780 nm, an energy per pulse of 800 μ J, and repetition rates of 1 kHz or 125 Hz. A 45° incidence angle was maintained, with *p* polarization. The laser was focused down on the target surface achieving area power densities which ranged between 2×10^{12} and 3.3×10^{13} W/cm². The target to substrate distance was 5 cm and the incident-beam power was varied using an attenuator. The ablation target manipulator can hold up to four different 1 in. targets with full individual rotation and selective clocking motion. Rotation and rastering allow for uniform removal of target material with successive laser shots, avoiding damage of the target by the beam.²² Two different sets of samples, with 1 min Co deposition for all samples, were grown on Si substrates. The first set consisted of Co deposits grown with different laser powers for morphological studies with atomic force microscopy (AFM) and scanning electron microscopy (SEM). A second set consisted of Au/Co/Au trilayers deposited with different laser powers and deposition times for Co. This second set of samples was used for compositional [by energy-dispersive x-ray spectroscopy (EDS)], structural [by transmission electron microscopy (TEM)], and magnetic characterization. Au was evaporated from an effusion cell with the evaporation rate calibrated by determining the thickness of Au layers from x-ray reflectometry. The thickness of the Au buffer and capping layers was 16 and 6 nm, respectively. Such capping layer thickness was thin enough to be able to extract information from the material underneath by means of x-ray magnetic circular dichroism (XMCD) and had proven to be effective in protecting Co layers and nanodisc layers against oxidation in previous studies where the samples had been prepared by means of MBE, sputtering, and lithographic processes.^{34,35}

Si substrates were cut into 1×1 cm² squared pieces, cleaned ultrasonically in baths of trichloroethylene, acetone, and methanol, blow dried with nitrogen, and loaded into the deposition chamber. Prior to the growth of the different layers at room temperature (RT), the Si substrates were outgassed at 600 °C for 45 min. AFM and magnetic force microscopy (MFM) images were taken at ambient conditions using a NanotecTM microscope operating in a noncontact dynamic mode. For MFM, the tip was scanned at a constant lift height above the sample, typically 50 nm, and the phase shift, proportional to the force gradient, was measured. Silicon probes from NanosensorsTM with a cantilever force con-

stant of ~ 3 N/m, resonant frequency of ~ 70 kHz, and pyramidal tip with radius below 10 nm, were used. For MFM, the tips were coated with Cr/Co/Cr trilayers by electron-beam (e-beam) thermal evaporation, with the thickness of Co being varied between 20 and 60 nm in order to have a suitable magnetic moment for each experiment: too low a magnetic moment would produce a poor signal-to-noise ratio, while a too high moment would alter the magnetic state of the nanoparticle being analyzed. All the tips were magnetized along the pyramid axis.

Structural and chemical characterization was carried out by high-resolution transmission electron microscopy (HR-TEM) and electron energy-loss spectroscopy (EELS) in a Jeol 2010 FEG at 200 keV, with a GIF spectrometer. The composition of trilayers was determined by EDS. High-resolution morphology studies of the deposited nanoparticles were made in a field-emission SEM (FEI Nova Nanolab) with 3.5 nm resolution operating at 30 keV. Magnetization measurements were carried out by using an rf superconducting quantum interference device (SQUID) magnetometer from Quantum Design in applied fields of up to 20 kOe.

X-ray photoemission spectroscopy (XPS) measurements have been performed in a separate UHV chamber with a base pressure of 1×10^{-10} mbar. The angle between the hemispherical analyzer (Specs-PHOIBOS100) and the plane of the surface was kept at 60° and used the Mg $K\alpha$ line (1253.6 eV) for the x-ray radiation. XPS spectra were recorded after removing a 15-Å-thick layer of materials from the samples by mild Ar^+ sputtering at 1 keV with a known ion sputtering rate.³⁶ The ion sputtering was performed in order to remove the contaminants from the surface and to enhance the photoemitted signal from the Co $2p$ core levels. It should be noted that the removal of the surface layers of the sample is very appropriate for the further quantification of oxygen in the sample since it removes the oxygen atoms adsorbed from the air. The spectra were recorded with photon energy steps of 0.1 and 15 eV pass energy. Prior to the data analysis, the contributions of the Mg $K\alpha$ satellite lines were subtracted and the spectra were subjected to a Shirley-type background subtraction formalism.

The magnetic characterization of the samples using XMCD was performed by measuring the x-ray-absorption spectra (XAS) in total electron-yield mode on beamline 5U.1 of the Synchrotron Radiation Source (SRS) at Daresbury, UK. All XMCD measurements were performed at room temperature using the portable octupole magnet system (POMS), which can reach magnetic fields up to ± 0.9 T. With the POMS the magnetic field can be applied in any direction, although the results presented here were measured by applying the magnetic field along the incident x-ray beam. The direction of the applied magnetic field was reversed at each photon energy point of the spectrum. The degree of polarization was 53.4% as deduced from a reference sample with the configuration 30 Å Al/40 Å Co/ $\text{Al}_2\text{O}_3(0001)$. Incomplete magnetization saturation was corrected with the values of the saturation magnetic field extracted from the Kerr hysteresis loops measured in polar geometry.

Due to the nanoparticle nature of the Co deposits, no direct determination of the deposition rate was possible, as, for example, with standard x-ray reflectivity (XRR) measure-

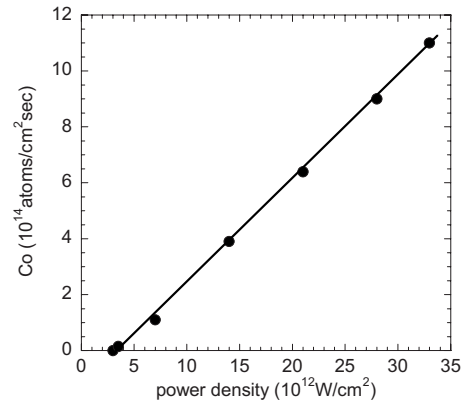


FIG. 1. Co deposition rate as a function of laser power density. A linear behavior is obtained, intersecting zero deposition at 3×10^{12} W/cm 2 , in good agreement with the experimentally observed ablation threshold.

ments of a continuous layer. The Co deposition rate was indirectly determined by EDS measurements from Au/Co/Au/glass trilayers where the Au thickness was known. Au grows with (111) texture on glass as well as on Co (Ref. 37) and therefore it is possible to know the number of Au atoms per unit area. Since EDS measurements provide the relative number of Co atoms per Au atom, it is therefore possible to extract the number of Co atoms per unit area and as a consequence the Co deposition rate. In Fig. 1 we show the Co deposition rates obtained as a function of laser fluence. As can be seen, the rate exhibits a linear behavior, starting from a threshold at around 3×10^{12} W/cm 2 . This agrees very well with the experimentally observed minimum power needed to ablate material off the target (ablation threshold).

III. CHARACTERIZATION

A. Morphology

1. Co/Si

To study the laser power dependence of the particle morphology and size distribution, Co was deposited onto oxidized Si substrates at different laser power densities ranging from 3.5×10^{12} to 3.3×10^{13} W/cm 2 with 1 min deposition time in all cases. The results are shown in Fig. 2, where SEM images are presented for representative powers. Ellipsoidal nanoparticles, close to circular and of different diameters, are observable in all cases, with a larger amount of nanoparticles for increasing power. For 3.5×10^{12} W/cm 2 , just above the ablation threshold, there is already a wide distribution of particle diameters from around 70 nm down to the resolution limit of the microscope. As expected, the number of particles increases with laser power, all the way up to the maximum power density of 3.3×10^{13} W/cm 2 . In some cases, and for all laser powers, an annular structure is observed, indicative of a disklike or ringlike geometry rather than spherically shaped particles. Coalescence, especially for high powers, and perimeter decoration with smaller particles are also observed.

In order to obtain the particle size distribution vs laser power, we have performed statistics from different SEM im-

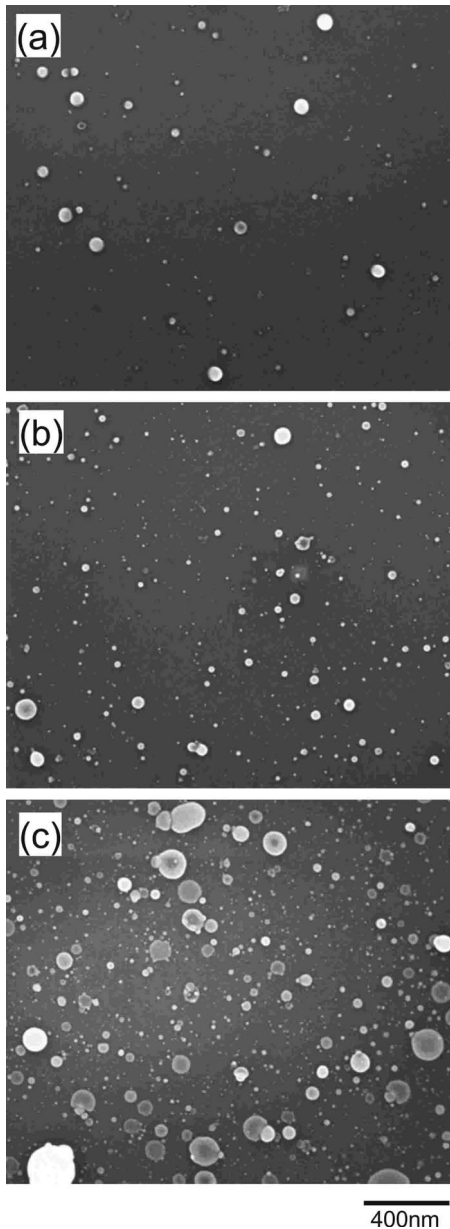


FIG. 2. SEM images obtained for 1 min depositions of Co on Si at (a) 3.5×10^{12} W/cm², (b) 7×10^{12} W/cm², and (c) 2.8×10^{13} W/cm².

ages for all the laser powers, measuring the in-plane diameter. For higher laser power density, the observed maximum disk size increases, as shown in Fig. 3; for instance, the largest in-plane diameter found for 2.8×10^{13} W/cm² is 240 nm, 150 nm for 7×10^{12} W/cm², and 120 nm for 3.5×10^{12} W/cm². We also performed a similar statistical on AFM images of selected samples, yielding similar results except for the small diameter range, where the convolution effect with the AFM tip radius (~ 10 nm) affects the results. Although previous studies reported a characteristic particle diameter between 5 and 10 nm,^{13,24,25} our experimental data do not show a peak in the particle size distribution, but rather a continuous increase in the number of particles as their size is reduced until the resolution of the available images prohibits determining such statistics. Another factor playing a

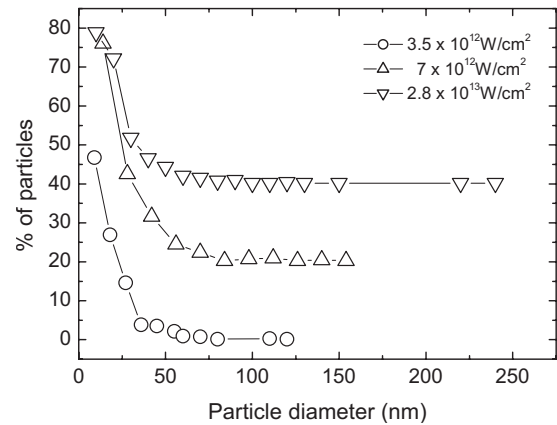


FIG. 3. In-plane particle diameter distribution (normalized to the total amount of computed particles) extracted from representative SEM images taken for Co/Si samples deposited at three selected laser power densities. The distributions for 7×10^{12} and 2.8×10^{13} W/cm² have been vertically displaced by 20% and 40%, respectively, for clarity. For all densities, the number of particles increases gradually and continuously as the particle diameter is reduced. However, for higher power density larger nanostructures can be produced.

role could be the different optical, thermal, and rheological properties of the materials studied in this work compared with Co, being therefore likely to produce slightly different particle size distributions.

Although the apparent morphology of the deposited particles in the SEM images is generally circular, some nanostructures appear as nanorings and nanoseries as can be observed in Fig. 4. Since SEM is not very sensitive in depth, AFM images were taken in representative samples to investigate the actual shape of the deposited Co nanostructures. The AFM measurements revealed that the shape of the particles referred to as circular is actually oblate semispherical, with an aspect ratio (diameter of the base divided by the polar axis) between 3 and 15. The shape of the nanorings is similar to that of a volcano, whereas the nanoseries are circular structures surrounded by small particles in the form of droplets. These findings indicate that the liquid state plays an important role in the formation of at least some of the nanostructures. One possibility is that Co clusters arrive in liquid state at the substrate and exhibit a plastic behavior after impinging on the surface, freezing into the finally observed morphology. Another interpretation is that Co clusters hit the substrate in a solid state, and then melt due to the energy delivered in the collision, turning back to a solid after thermalization with the substrate.

2. Au/Co/Au/Si

For further magnetic characterization, Au/Co/Au trilayered structures grown on Si and with 1 min deposition of Co at different power densities were used. AFM images for the different samples are shown in Fig. 5.

As can be observed, the Au layer does not completely cover the Si substrate, but deposits as Au aggregates, roughly 100 nm in lateral dimension and close to each other. For Co

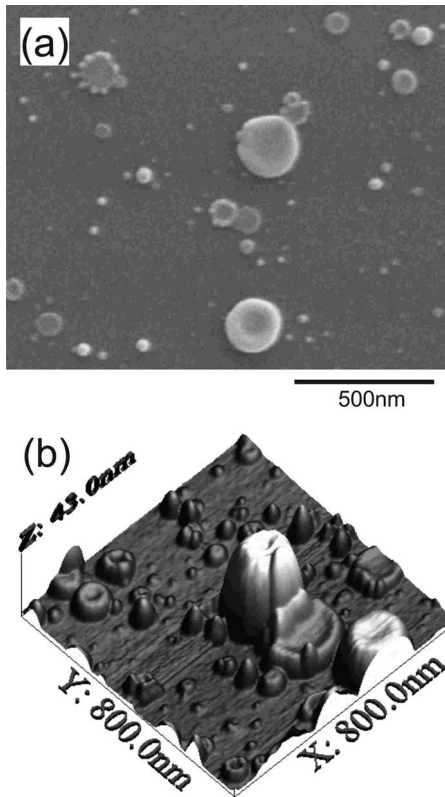


FIG. 4. Images for selected Co/Si samples grown at 3.3×10^{13} W/cm² exhibiting shapes different from disks. Both nanorings and nanorosaries are observed. (a) SEM image. (b) AFM image.

deposited at 2×10^{12} W/cm², just below the determined ablation threshold, the morphology is almost indistinguishable from the bare Au layer, implying a very small Co particle size or simply the absence of Co in the layer (keeping in mind the 3×10^{12} W/cm² ablation threshold). No particle is observed by AFM and only one feature, which could be a Co particle, was observed in the SEM measurements. This indicates that the amount of Co ablated at this subthreshold laser power is very small.

At 3.5×10^{12} W/cm² power density, a few features on top of the Au layer are observable, both by AFM and SEM. Diameters and heights vary, but range roughly between 50 and 200 nm in diameter and 10–20 nm in height. This implies that the particles are actually disks with an aspect ratio that might actually depend on the aggregate size. For example, the smaller Co particles are 50 nm in diameter and 20 nm high, while others, although a minority in number, have a distinctive ring shape, with a central dip that may or may not leave the central part of the ring covered with Co. This is more noticeable for 7×10^{12} W/cm², where Co rings of two types, either with the central part covered or uncovered with Co, are observed. For 1.4×10^{13} W/cm², a larger population of Co particles is clearly observed, even though in this specific case no ringlike structures were detected. The same trend, with increasing amount of Co particles as a function of the laser power, is observed for the 3.3×10^{13} W/cm² samples.

Finally, regarding the structural characterization, TEM analysis was carried out on selected Au/Co/Au/Si samples, in

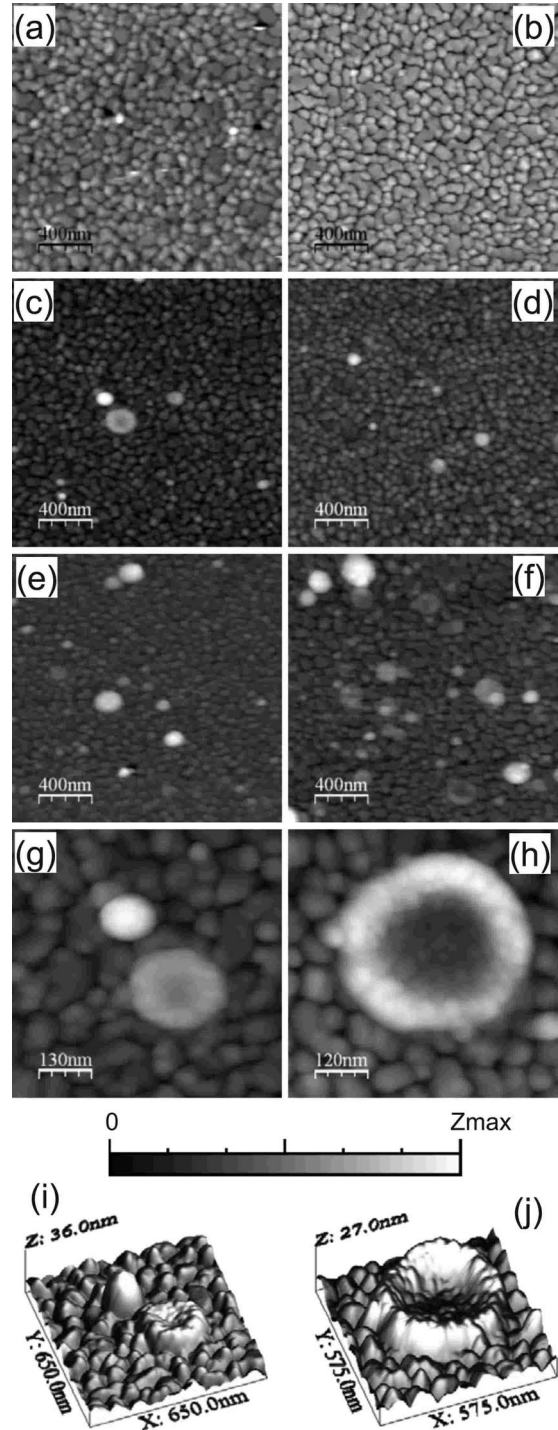


FIG. 5. AFM images for Au/1 min Co/Au/Si samples. [(a)–(f)] Representative top views of samples deposited at 0, 2, 3.5, 7, 14, and 33×10^{12} W/cm², respectively. Selected images of specific particle shapes: [(g) and (h)] top views and [(i) and (j)] 3D images of nanostructures obtained at 3.5 and 7×10^{12} W/cm², respectively.

both plan-view and cross-section geometries [Figs. 6(a) and 6(b), respectively]. Both plan view and cross sections confirmed the discontinuity of the Au layer as seen by SEM. Samples in cross-section geometry were observed by EELS and HRTEM. An example corresponding to a sample grown

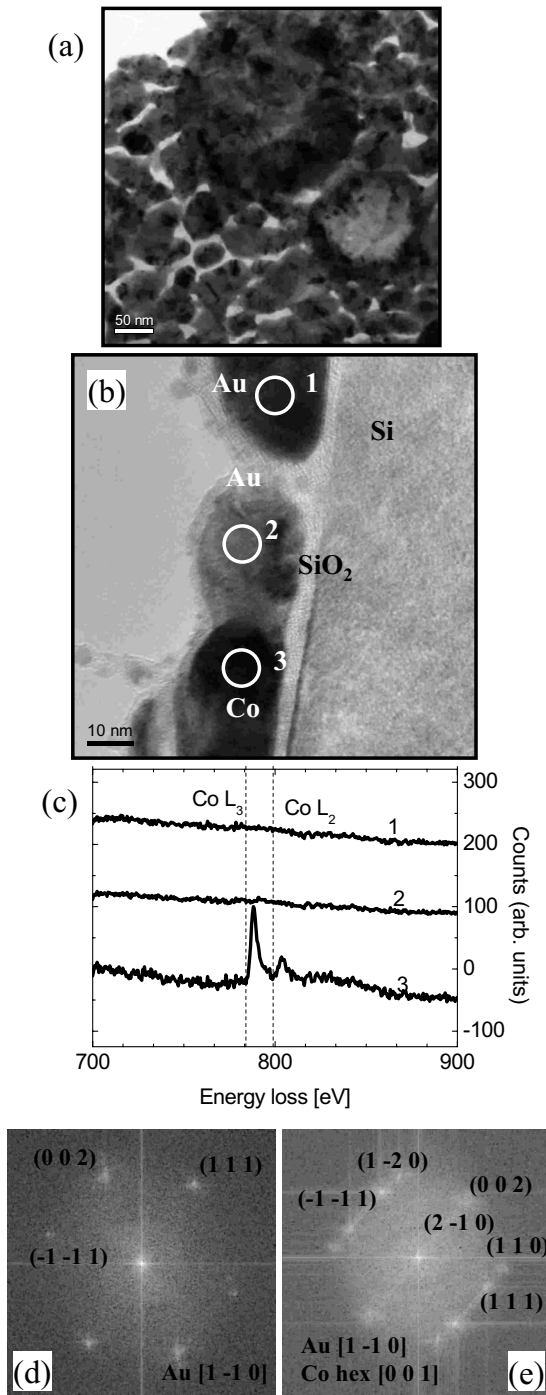


FIG. 6. (a) Plan-view image of 3.3×10^{13} W/cm² sample, where Co islands can be observed. (b) Cross-section image of 1.4×10^{13} W/cm² sample, displaying a Co particle with a diameter of about 19 nm. (c) EEL spectra in the indicated zones in (b). (d) FFT of the gold particle in (b). (e) FFT of the Co particle in (b), where $(-1-11)$, (002) , and (111) reflections correspond to Au, and $(1-20)$, $(2-10)$, and (110) reflections correspond to hcp Co.

at 1.4×10^{13} W/cm² is shown in Fig. 6(b). EELS allowed distinguishing Co particles from the Au buffer and capping discontinuous layers [Fig. 6(c)]. Within the experimental limits of the technique, no evidence of a continuous Co layer connecting the Co nanoparticles could be observed. Power

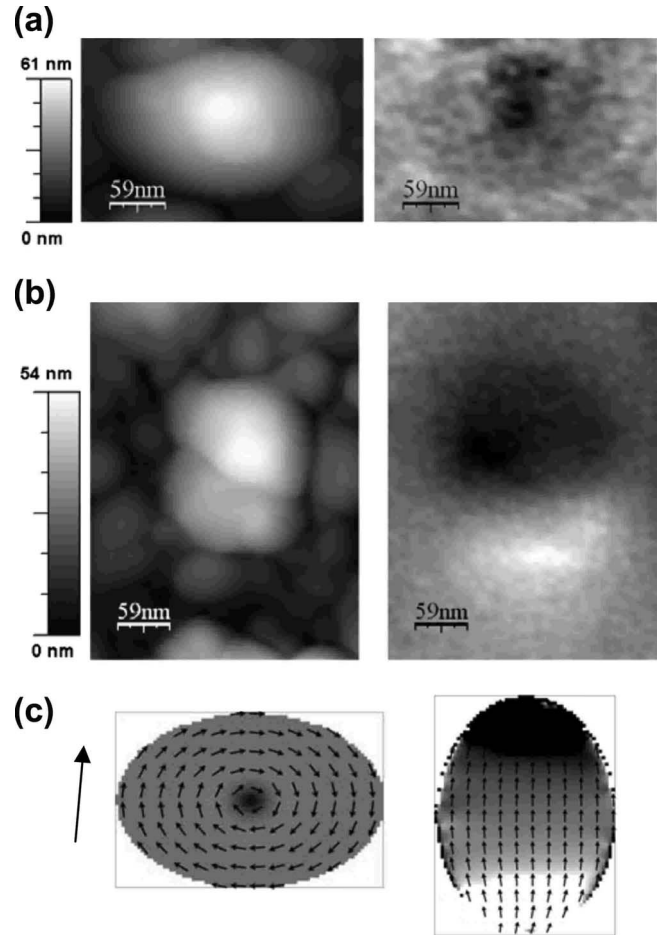


FIG. 7. [(a) and (b)] MFM images (left: topography; right: magnetic signal) for two nanoparticles exhibiting: (a) vortex state, (b) single-domain state. (c) Micromagnetic simulations for hcp nanoparticles in the remanent state after assuming a field of 2 T along the direction indicated by an arrow exhibiting: (left) vortex state, (right) single-domain state; the gray scale represents the magnetic charges (poles).

spectra calculated by fast Fourier transform of HRTEM images also helped to identify and discriminate between Au and Co particles, the latter exhibiting an hcp structure [Figs. 6(d) and 6(e)]. Selected area diffraction patterns (not shown) evidenced also the presence of bright spots corresponding to CoO.

B. Magnetic properties

1. MFM and micromagnetic simulations

Due to the disklike shape of the deposits, it is of interest to study the local magnetization distribution within the Co entities. This was carried out by MFM in the remanent state after in-plane saturation, although only the larger nanoparticles could be observed since the lateral resolution of MFM at ambient conditions is about 50 nm. Two magnetic states have been found: the vortex (V) state as shown in Fig. 7(a), and the single-domain (SD) state as shown in Fig. 7(b). The V state consists of a curling spin configuration that develops a small core at the center with out-of-plane magnetization.

Such a core creates a magnetic pole at the surface, and as MFM is a field line mapping microscopy, it gives rise to a monopolar contrast in the magnetic signal,³⁸ i.e., the black dot at the center in the right image of Fig. 7(a). Taking into account the gray scale used in the experiments (black for attractive interaction, white for repulsive interaction), we can state that the V core magnetization is parallel to the tip magnetization. Antiparallel tip and core magnetizations would have produced a white dot at the V core. On the other hand, in the SD state the magnetic moments are mainly aligned along one particular in-plane direction and consequently create opposite magnetic poles at the sample ends. Such poles give rise to a dipolar contrast in the magnetic signal, i.e., black contrast at one end and white contrast at the opposite end of the nanoparticle, as shown in the right image of Fig. 7(b). The fact that the magnetic image is asymmetrical (the black end appears to be larger than the white one) has been described elsewhere and is due to the tip-sample interaction in the MFM experiment.³⁹

It is worth mentioning that in our experiments the SD state was preferentially found in the ellipsoidal particles with the long axis not far from the direction of the initially applied magnetic field. Finally, in order to shed light on these observations, we performed micromagnetic simulations using the OOMMF code⁴⁰ (parameters: exchange constant $A=3 \times 10^{-11}$ J/m, saturation magnetization $M_s=1.4 \times 10^6$ A/m, cell size=3 nm, i.e., below the exchange length for Co). It is well known that the magnetically stable states exhibited by Co dots depend on the element size (lateral dimensions and thickness), crystalline anisotropy, and magnetic history.^{41,42} The calculations have been performed considering two ellipsoidal nanoparticles with similar thickness (17 nm) but different axes along the x and y directions: (252 nm, 165 nm) for one of them (named particle A) and (168 nm, 225 nm) for the other (particle B), mimicking the lateral dimensions of the particles shown in Figs. 7(a) and 7(b), respectively. In order to reproduce the MFM measurements, an initial in-plane applied field of 2 T deviating 5° from the y axis has been considered, and the analyzed configurations correspond to the remanent states once the field is removed. The results shown in Fig. 7(c) were obtained assuming that the particles have hcp crystalline structure with magnetocrystalline easy axis always oriented along the surface normal, the anisotropy constant being $K=5.2 \times 10^5$ J/m³: particle A (on the left) exhibits a remanent V state similar to that experimentally observed in Fig. 7(a), whereas particle B (on the right) exhibits a remanent SD state as that experimentally determined, as shown in Fig. 7(b). However, if an initial out-of-plane field is used instead of in-plane field, which was considered to reproduce the MFM measurements, both particles reach a remanent V state. Comparing the total energy density E of the two stable magnetic states obtained for particle B, i.e., SD for the initial in-plane field ($E=6.1 \times 10^5$ J/m³) or V for the previous out-of-plane saturation ($E=5.5 \times 10^5$ J/m³), we can deduce that the V configuration is the ground state.

From a magnetic viewpoint, it is crucial to know whether or not the magnetocrystalline easy axis has an important component along the surface normal. In order to qualitatively study the influence of this parameter on the development of the remanent state configurations after in-plane saturation,

simulations varying the anisotropy constant in the perpendicular direction were performed. When the anisotropy is halved, the obtained configurations are similar to that shown in Fig. 7(c), i.e., particle A exhibits a V state whereas particle B (on the right) exhibits an SD state. However, when the anisotropy constant is reduced by a factor of five, the remanent state for particle A is no longer a V but a SD state with the magnetic moments mainly aligned along the long axis. This means that for the elliptical particles considered here, the development of a remanent V state (where the moments at the V core are pointing out of the plane) after in-plane saturation requires a global (effective) out-of-plane magnetocrystalline anisotropy high enough to overcome the shape influence. In other words, the experimental observation by MFM of remanent V states points to the existence of a certain texture in the crystalline structure of the hcp particles with the c axis (easy-magnetization axis) pointing out of the plane. Although the TEM measurements indicate that the nanoparticles exhibit hcp structure, the number of Co particles examined does not give significant statistics about the orientation of the easy axis.

Finally, in order to obtain an overview of the influence of the magnetocrystalline anisotropy in the development of the magnetically stable states after in-plane saturation, we performed simulations considering circular particles (in order to remove the additional shape anisotropy) with fixed thickness (17 nm) and different diameters. In clear contrast to the ground-state case, in which the presence of perpendicular magnetic anisotropy has little influence on the critical size for the transition from the SD ground state toward the V ground state,⁴³ we observed a strong influence of the anisotropy constant in the magnetic state obtained after in-plane saturation. In particular, for perpendicular anisotropy with $K=5.2 \times 10^5$ J/m³ the critical diameter is 125 nm (the remanent configuration for diameter below or above such value is the SD state or the V state, respectively), whereas for $K=2.6 \times 10^5$ J/m³ the critical diameter increases to 362 nm.

2. SQUID

For the determination of the saturation magnetization and the possible presence of superparamagnetic nanoparticles, in-plane hysteresis loops at low temperature (10 K) and RT were measured in all samples using a SQUID magnetometer (Fig. 8, with m being the magnetic moment), confirming their ferromagnetic nature up to RT. All loops exhibit a similar shape, with no observable variation in m_s (magnetic moment at saturation) upon reduction in temperature. The fact that the remanent moment is larger at low temperature can be easily understood by considering the configurations discussed in the previous section. Such a remanent moment is due to the particles that exhibit SD configurations (the V state does not produce a net in-plane moment at zero field), and upon increasing temperature the thermal energy induces the transition to the ground state (i.e., the V state) of some of them, thereby reducing the remanent moment. In Figs. 9(a) and 9(b), RT in-plane and out-of-plane hysteresis loops are shown for the 7×10^{12} and 2.1×10^{13} W/cm² samples, all of them confirming a clear in-plane magnetic anisotropy most likely produced by the morphology of the particles. The rem-

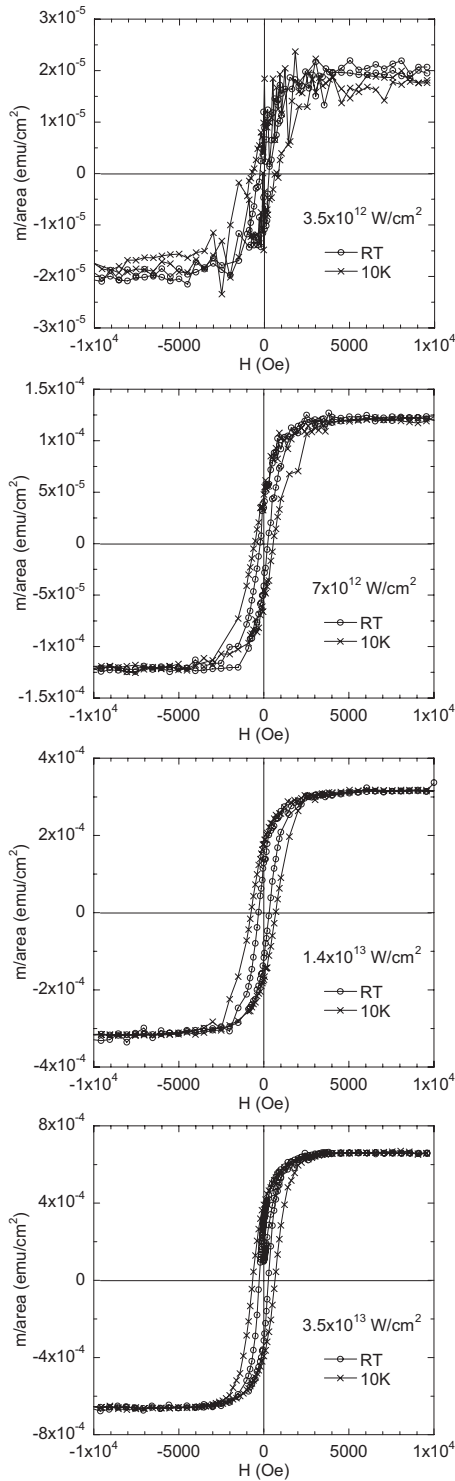


FIG. 8. Magnetic hysteresis loops vs laser power obtained from in-plane SQUID measurements at room temperature and 10 K, showing their ferromagnetic nature up to RT.

alent moment for the out-of-plane loops can be mainly attributed to the spins at the core of the particles exhibiting V states. It is worth noticing that, in spite of the minute size of the smallest detected particles, no evidence of superparamagnetism is observed in the loops. Since the measured magnetic moment could be due mainly to the larger particles, several

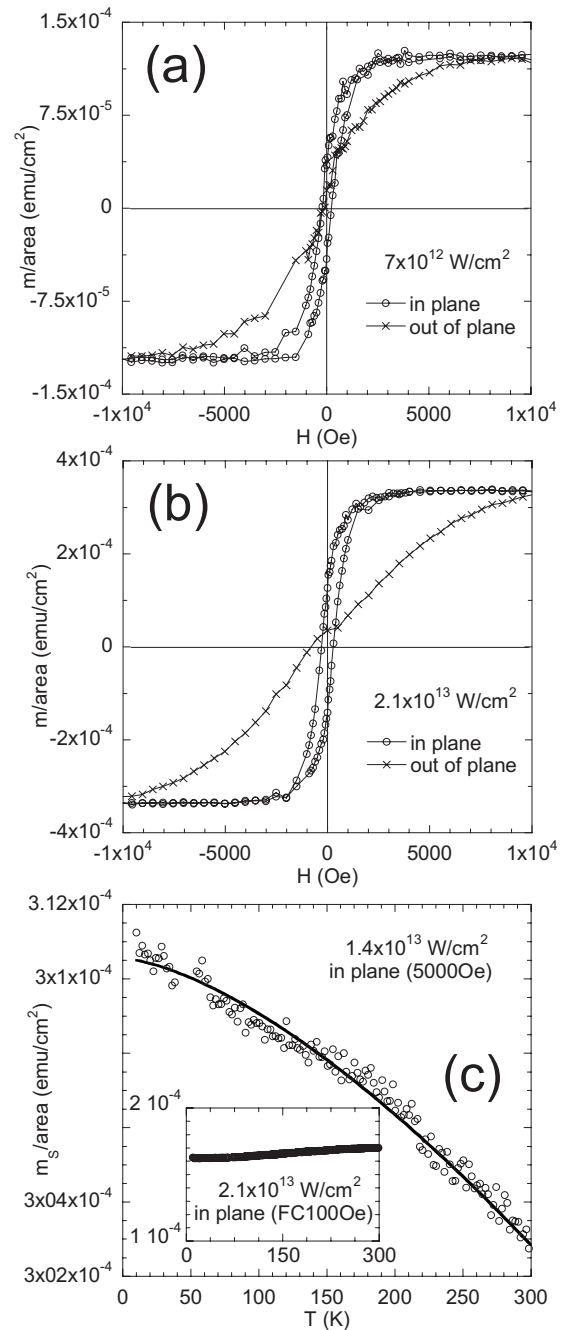


FIG. 9. [(a) and (b)] In-plane and out-of-plane magnetic hysteresis loops at RT of (a) 7×10^{12} W/cm² and (b) 2.1×10^{13} W/cm² samples, showing an in-plane magnetic anisotropy. (c) m_s vs T for the 1.4×10^{13} W/cm² sample, line fit corresponding to Bloch $T^{3/2}$ law. The inset shows the FC magnetization curve at 100 Oe for the 2.1×10^{13} W/cm² sample.

temperature-dependent measurements have been performed to elucidate this point.

The m_s vs T curve is presented in Fig. 9(c) for the 1.4×10^{13} W/cm² sample, showing the Bloch $T^{3/2}$ law, with a B factor of 4.8×10^{-6} and with only 3% reduction in m_s when increasing the sample temperature from 10 K to RT. This indicates that most of the Co aggregates are already ferromagnetic at RT, and do not exhibit any relevant amount of superparamagnetic particles and/or isolated atoms in the

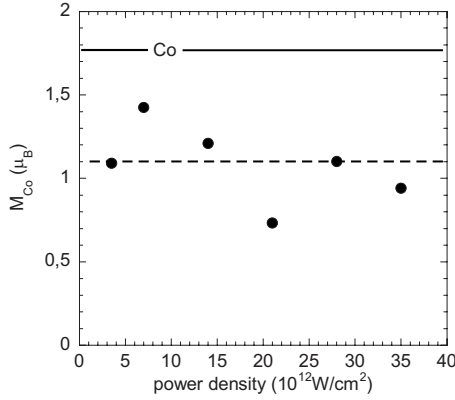


FIG. 10. The Co magnetic moment per atom estimated from the Co deposition rate, showing a low magnetic moment of $1.1\mu_B$, compared to $1.77\mu_B$ for bulk Co.

films. Another indication that points to the same conclusion is given in the inset of Fig. 9(c): the field-cooled (FC) magnetization curve measured with a magnetic field of 100 Oe for the 2.1×10^{13} W/cm² sample is displayed. There is no evidence for increased magnetization values at low temperatures as those expected for superparamagnetic particles and isolated ions. Only a monotonous increase in m with increasing temperature is obtained, which can be understood if the assembly of ferromagnetic nanostructures was away from its equilibrium state.

Using the estimated Co deposition rate (Fig. 1) it is possible to determine the amount of Co in the structures and therefore the magnetic moment per Co atom as a function of the laser power (Fig. 10). In spite of the low accuracy due to both the small amount of Co in the samples and the errors in the estimation of their volume, an average value of around $1.1 \mu_B$ is obtained, which represents about a 40% reduction with respect to bulk Co. As will be shown in the following sections, this significant reduction in the magnetic moment can be attributed to a partial oxidation of Co.

3. XMCD

An independent determination of the Co magnetic moment value was carried out by XMCD measurements. Figure 11(a) displays the Co $L_{2,3}$ XAS measured for opposite magnetization directions (μ^+ and μ^-) on the sample with 1.4×10^{13} W/cm². A clear dependence of the absorption intensity can be observed on the magnetic polarization of the sample, as expected for ferromagnetic systems. The difference in the absorption intensities gives a direct measure of the magnetic nature of the sample at room temperature. Figure 11(b) shows the corresponding sum and difference spectra, the latter referred to as the XMCD signal, together with their corresponding integration curves obtained from a standard analysis procedure, where a double steplike function set to the peak positions of the L_3 and L_2 maxima was subtracted from the isotropic absorption spectra in order to remove the L_3 and L_2 edge jumps. The Co magnetic moments were extracted using the sum rules.^{44–46} Here we neglected in the analysis the so-called saturation effects,^{47,48} mainly because of the small thicknesses of the structures studied here. The

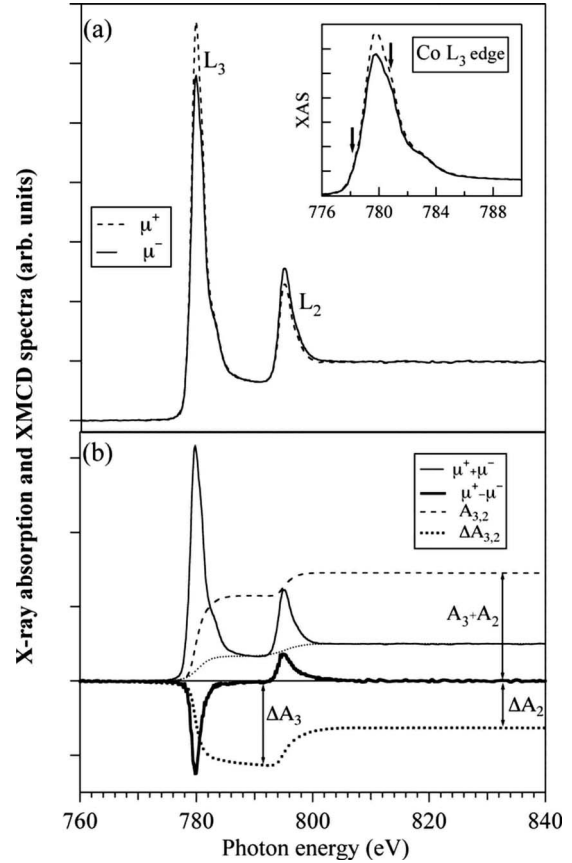


FIG. 11. (a) Normalized XAS data measured on the sample grown at 1.4×10^{13} W/cm² and for opposite magnetizations. (b) Normalized summed XAS spectra, the corresponding difference spectrum (XMCD spectrum), and their integration curves. The dotted line is the two-step-like function for the edge-jump removal before the integration. The ΔA_3 , ΔA_2 , and A_3+A_2 values are the three different integrals required for the sum-rule analysis.

relation between the orbital moment μ_L , spin moment μ_S , magnetic-dipole term T_z , and number of $3d$ holes, n_h (all per atom), in the ground state and the intensities, ΔA_3 , ΔA_2 , and A_3+A_2 of the difference $\Delta A_{2,3}$ and sum $A_{2,3}$ spectra integrated over the corresponding $L_{2,3}$ edges is given by^{44–46}

$$\frac{\mu_L}{n_h} = -\frac{4}{3} \frac{\Delta A_3 + \Delta A_2}{A_3 + A_2}, \quad (1)$$

$$\frac{\mu_S + 7T_z}{n_h} = -2 \frac{\Delta A_3 - 2\Delta A_2}{A_3 + A_2}. \quad (2)$$

Theoretical results⁴⁹ indicate that $7T_z/\mu_S < 0.1$, hence giving an error in μ_S of up to 10% by neglecting T_z .

The extracted values for the Co magnetic moments are given in Table I together with values for bulk Co taken from Ref. 46 for comparison. The values have been obtained assuming the standard value of $n_h = 2.49$ for cobalt in the hcp structure.⁴⁶ We assume that the mixing of the $2p_{3/2}$ and $2p_{1/2}$ states as well as the contribution of the sp valence band to the $L_{2,3}$ XMCD signal can be neglected and that $T_z = 0$.

TABLE I. Co spin and orbital atomic magnetic moments, μ_L and μ_S , their ratio and the total atomic magnetic moment μ_T (in μ_B) extracted from XMCD and XAS.

	μ_S (μ_B)	μ_L (μ_B)	μ_L/μ_S	μ_T (μ_B)
Bulk Co (Ref. 46)	1.62	0.154	0.095	1.77
Expt. (XMCD)	0.95(1)	0.190(2)	0.199(1)	1.14(1)

The total atomic magnetic moment obtained for this sample by XMCD is smaller than the $1.77\mu_B$ reported for bulk hcp cobalt.⁴⁶ The reason for this behavior could be the presence of oxidized Co. In fact, the appearance of multiplet structure [see inset of Fig. 11(a)] superimposed on the metallic Co spectrum⁵⁰ strongly suggests the formation of cobalt oxide. The extra peaks in the absorption spectrum, indicated by the arrows in the inset, agree very well with those observed in the measured spectrum of CoO as well as with the calculated multiplet.⁵¹ Antiferromagnetic CoO does not contribute to the dichroic signal so that the XMCD in Fig. 11(b) only displays the metallic line shape. It is important to note here that the application of the sum rules gives the average values of the magnetic moments per atom since the denominators in Eqs. (1) and (2) include all Co atoms, i.e., metallic as well as oxidized Co atoms. Actually, the total atomic magnetic moment measured corresponded to 64% of that from bulk cobalt, in very good agreement with the values obtained with SQUID. Imperfections in the capping layer could lead to two different systems. On one hand, the system could be formed by a portion of the particles completely converted into oxide, while keeping the rest in the metallic form. On the other hand, the system could be formed by Co disks with a metallic core surrounded by an oxide shell.

The analysis of the μ_L and μ_S contributions reflected significant differences when compared to bulk cobalt. There is a strong decrease in the spin magnetic moment. Some authors reported that the high-spin orbit of gold in very small particles induces an enhancement of the μ_L/μ_S ratio of the metal transition at the interface.^{52,53} However, in the systems studied here, this effect is likely to be very attenuated as a consequence of the bigger particle size and, thus, the observed decrease in the μ_S would more likely be indicative of an oxidized interface. Incomplete oxidized interfaces would lead to an increase in the effective Co spin polarization (μ_S/n_h). However, once the interface is saturated with Co-O bonds, the effective Co spin polarization is reduced as more Co atoms are incorporated into the cobalt oxide.⁵⁴ As a consequence, the small effective Co spin polarization measured by XMCD suggests the formation of a continuous paramagnetic cobalt oxide at the Co interface.

On the other hand, an increase in the orbital magnetic moment was measured compared to bulk cobalt. For the nanoparticles the orbital to spin magnetic-moment ratio is roughly doubled (see Table I). This result is in agreement with earlier observations on Co clusters that found a strong variation in the atomic magnetic moment with cluster size.^{55–57} The orbital magnetic moment is directly linked to the electronic structure and its size depends on the 3d hy-

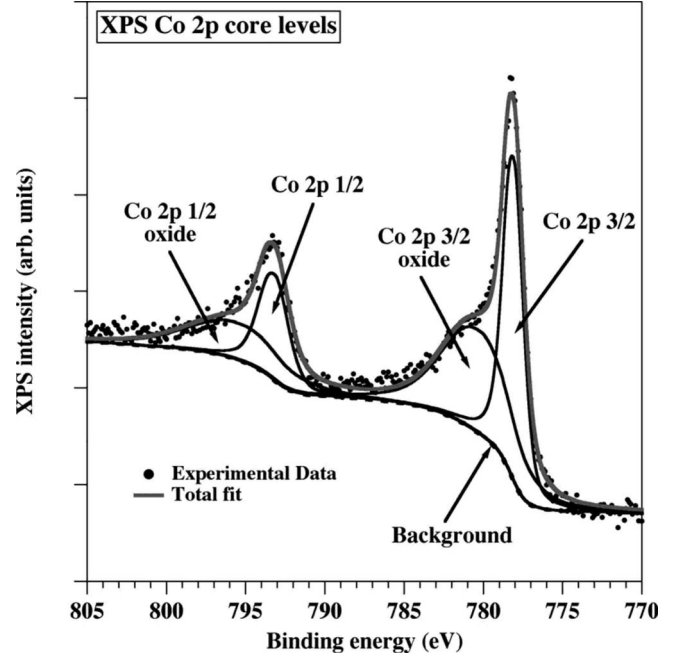


FIG. 12. Representative Co 2p core-level XPS spectrum of the sample grown at 1.4×10^{13} W/cm².

bridization and localization. The increase is generally ascribed to electron localization at the perimeter of the clusters. The orbital magnetic moment is very sensitive to small changes in the surface structure⁵⁸ and strongly depends on the geometry of the particle.⁵⁹ An enhancement also occurs if the Co is at a site where its orbital moment is not quenched by the crystal field, which for instance is the case for Co²⁺ in octahedral symmetry, such as for CoO.⁵¹ However, to observe this in the XMCD spectrum would require magnetic moments that are not antiferromagnetically coupled.

C. Chemical characterization: XPS

A very likely reason for the observed reduced magnetic moment with respect to that of the bulk is the oxidation of Co due to the discontinuous nature of the Au capping layers, and therefore their effective role against oxidation of Co was directly checked by XPS. A representative Co 2p core-level spectrum is displayed in Fig. 12. Together with the experimental data points, the background and the fitting curves are displayed. As can be observed, the experimental spectrum has been fitted with four components that correspond to the Co 2p_{1/2} and 2p_{3/2} core levels and their corresponding Co oxide components. From the integrated area of the components of the Co core levels and from the photoemission peaks of gold and oxygen (not shown here) and taking into account the corresponding photoemission cross sections, the proportion of Co oxide compared to the metallic Co was found to be $\sim 50\%$. More precisely, the chemical composition extracted from the XPS data is presented in Table II. The proportion of cobalt oxides and pure metallic cobalt found by XPS ($\sim 50\%$) is in good agreement with the reduction in the total atomic magnetic moment determined by XMCD ($\sim 36\%$ as compared to the atomic magnetic moment of pure

TABLE II. Chemical composition as deduced from the XPS spectra (at. %) from representative sample.

Au $4f_{5/2}$	Au $4f_{7/2}$	O $1s$	Co $2p_{1/2}$	Co $2p_{3/2}$	Co $2p_{1/2}$ oxide	Co $2p_{3/2}$ oxide
38	38	14	2.3	2.2	2.0	2.3
	76	14		4.5		4.3

metallic hcp Co) and SQUID ($\sim 40\%$ as compared to the atomic magnetic moment of pure metallic hcp Co). Such results clearly evidence the noncontinuity of the Au capping layer that induces a partial oxidation of the Co particles and in turn a reduction in the effective magnetic moment or magnetization.

Finally, it is worth mentioning that the existence of antiferromagnetic CoO usually induces unidirectional anisotropy in Co particle systems.⁶⁰ However, in our samples no shift in the hysteresis loops was observed at low temperature after field cooling, i.e., no exchange bias has been detected. It is very likely that the oxide is irregular and nonstoichiometric so the amount of antiferromagnetically ordered CoO is small and not able to produce exchange bias in the structures.

IV. DISCUSSION

Focusing first on the structural and morphological aspects, we systematically observed the formation of nanoparticles for all ablation powers above the ablation threshold, with an expected increase in the amount of deposited material with laser power. The particle size distribution does not show any noticeable differences when the laser power is increased, except for the observation of larger particles at larger powers. An important aspect is the absence of a characteristic peak for a specific particle size in the size distribution, with instead a continuously increasing number of particles for smaller diameters. This result, in contrast with the findings of other authors reporting a characteristic particle size, is in agreement with theoretical calculations that describe the ablation process as the emission of particles with a wide range of sizes, all the way down to the atomic size. Comparing with ablation techniques that use longer laser pulses, namely ns-PLD, where the presence of a high-pressure gas is necessary to obtain nanoparticles^{61,62} in our specific case nanoparticles are spontaneously obtained from the ablation process. The specific research reported here on the properties of nanoparticle systems provides a technically easier fabrication method of the desired structures. From morphological viewpoint, the obtained structures have systematically a disklike, instead of a spherelike shape, indicative of a liquid or at least plastic state of the particle impinging on the substrate, being deformed to this final shape. It is noticeable that, in addition to flat circular disks, in some cases other morphologies are also observed, such as rings and rosaries, which indicate a matter redistribution of the particle upon arrival at the substrate surface. This is of special interest in the case of magnetic materials since these shapes could potentially give rise to magnetic domain structures and anisotropies characteristic of these morphologies.

Two series of samples were grown, one directly on Si and the other one on a Au buffer layer on glass and subsequently covered with Au. The morphology of the Au buffer layer is not flat over large lateral dimensions, but is characterized by a granular nature due to the poor wetting of Au on glass. However, the final morphology of the particles in the Au/Co/Au structures is very similar to the Co/Si ones, indicating first that the Au capping layer must be substantially conformal with the Co structures underneath, reproducing their shapes; and second that the morphology of the Co structures does not strongly depend on whether these are deposited on Si (or actually its native oxide) or Au. This last point indicates that the final obtained morphologies are mainly determined by the state of the Co particles before arrival at the substrate, probably because these are in a liquid or at least a plastic state with a kinetic energy sufficient to give rise to particle deformation which produces the final observed shapes.

From magnetic properties viewpoint, only studied in the Au/Co/Au samples, we obtain a reduced magnetic moment per atom basically independent of the laser power. Within a rather wide range of laser power values, varying between 1.4 and $0.7\mu_B$ while the bulk value is $1.77\mu_B$. This result, confirmed by both SQUID and XMCD measurements, is due to the partial Co oxidation probably caused by a poor coverage of the Au capping layer. Both SQUID and XMCD are experimental techniques that probe large lateral areas, and therefore, even though the studied specimens present a wide range of particle sizes, the obtained results correspond to an average over the whole macroscopic system. In this respect, we have also applied a microscopical magnetic characterization technique such as MFM, which allows obtaining the domain structures of some of the observed morphologies, which are well reproduced by micromagnetic simulations. Moreover, these magnetic states can explain the macroscopic hysteresis loops measured by SQUID, both for in-plane and out-of-plane fields. On the other hand, we do not find any experimental evidence for superparamagnetism in the SQUID measurements. This is consistent with the smallest particle size detected in our SEM and AFM measurements, but it would be a surprising result if even smaller Co particles were present.

It is obvious that, for fs-PLD to become of relevance in a near future as a standard deposition technique for the fabrication of magnetic nanostructures, many aspects need to be improved, especially concerning the particle size distribution. However, we hope that this work will contribute to a deeper understanding of the morphological, structural, and magnetic properties of magnetic systems obtained by fs-PLD.

V. CONCLUSIONS

A comprehensive study covering structural, morphological, magnetic, and chemical characterization has been performed in Co structures grown by fs-PLD. The potential of fs-PLD as a powerful technique for the fabrication of magnetic entities of nanometric dimensions over large areas has been also demonstrated. Usually, nanodisks are obtained—their maximum size as well as the particle number increasing with the laser power density. Particles covered with Au were studied magnetically, showing no indication of superparamagnetism in spite of the small size of many of them. The measured magnetization exhibits some evolution with the laser power; however, the uncertainty in the determination of the actual amount of Co and the partial oxidation due to the discontinuous Au layer deposited on top make further studies necessary to clarify this issue.

Two remanent states have been identified after in-plane saturation: vortex and single-domain states, with the global magnetism dominated by the monodomain particles. We want to emphasize that the noncontinuity of the Au capping layer that induces the partial oxidation of the Co particles is an important aspect of the findings and contrasts with the behavior obtained in structures fabricated by, for example, MBE or sputtering. It highlights the different nature of the fs-PLD grown Co nanoentities with respect to those obtained

by other physical vacuum deposition methods.

The magnetic properties found in these structures are consistent with the presence of such a partial oxidation. The reduction in the effective magnetic volume explains the reduction in magnetic moment with respect to the bulk. The nonevidence of superparamagnetism is not only due to the disklike shape of the particles that increases the effective anisotropy, but is probably also related to the fact that the smallest particles are completely oxidized and consequently became nonmagnetic.

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